



SciDAC

Scientific Discovery through Advanced Computing



**AFOSR**

AIR FORCE OFFICE OF SCIENTIFIC RESEARCH



# Numerical Experiments with Shock Turbulence Interaction

Sanjiva K. Lele

Johan Larsson, Soshi Kawai

**Stanford University**

*Celebrating 50 Years of Turbulence and Acoustics Research*  
Symposium Honoring Prof. G. Comte Bellot, ECL, Lyon, France

**October 28-29, 2009**

Supported by the Department of Energy under the SciDAC -II program, and AFOSR-MURI program  
Computations performed at the National Energy Research Scientific Computing Center under the ERCAP program,  
and at Argonne National Laboratory under the INCITE program.

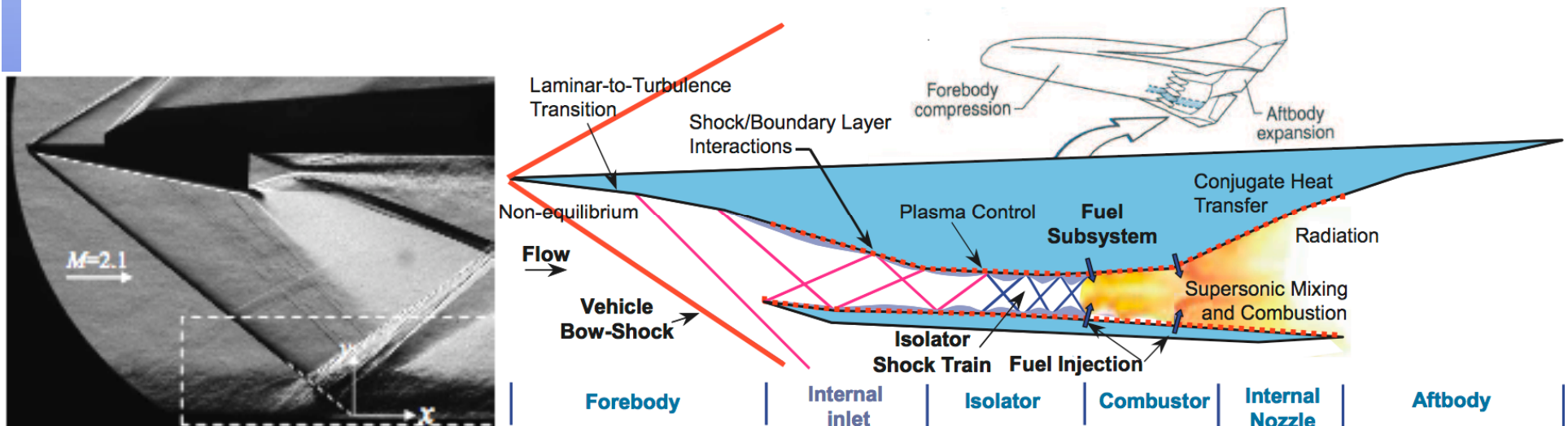
# Motivation (1)



Shock-turbulence interaction is central to high-speed propulsion systems yet its current models are quite inaccurate

Examples: scramjet engine; supersonic inlets

- Control of shockwave-turbulent boundary interactions in supersonic aircraft inlets (bleed system) causes a 15-20% performance penalty;
- Stability of the shock-train system limits the operability of a scramjet engine

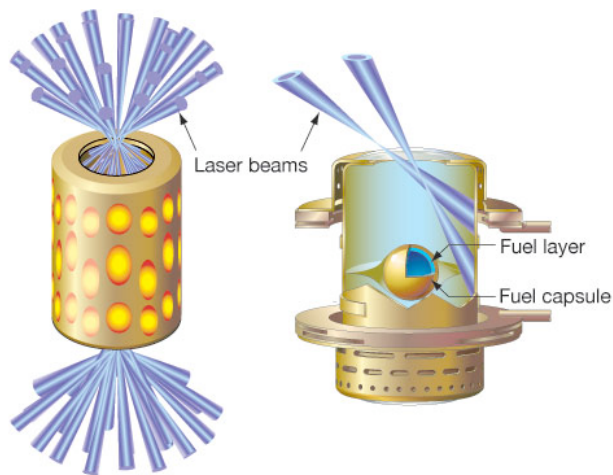


Shock-Boundary layer interaction

Scramjet engine

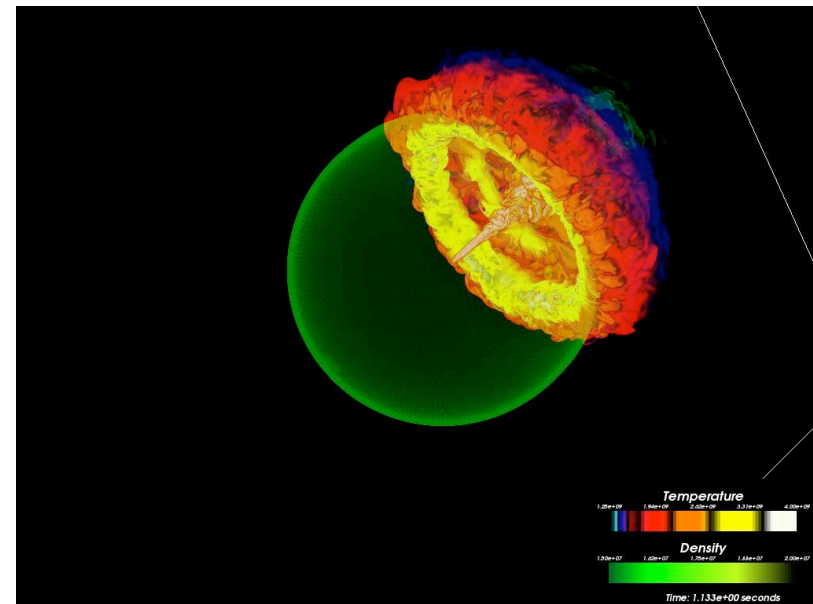
## Motivation (2)

- Interactions between strong shocks and interfaces lead to flow instabilities, which drive turbulent mixing
- Accurate prediction of shock-induced turbulent mixing remains an open problem
  - Examples: Rayleigh-Taylor mixing, Richtmyer-Meshkov mixing
  - Applications: inertial confinement fusion, supernova explosions



Source: [lasers.llnl.gov/programs/nic/icf/](http://lasers.llnl.gov/programs/nic/icf/)

Inertial confinement fusion by shock-induced implosion of a deuterium/tritium capsule



Source: [flash.uchicago.edu](http://flash.uchicago.edu)

Supernova explosion

# The key challenge:



- Simulation of broadband **turbulence requires** numerical methods with **minimum numerical dissipation**
  - Capture the widest range of resolvable scales on a given grid
  - Spectral methods; high-order methods
- Simulation of flows with **shocks requires numerical dissipation** for shock capturing
- The key challenge:  
  
How to design a numerical method that handles these **conflicting requirements**?



# Assessment of Numerical Algorithms: (for shock-turbulence interaction simulations)

- Broad range of high-resolution algorithms
  - High-order WENO (7th order)
  - High-order compact differences with selective artificial viscosity (STAN; STAN-I)  
Cook & Cabot, 2007, and modified Kawai & Lele, 2008;  
Mani, Larsson & Moin, 2009, Bhagatwala & Lele, 2009
  - Hybrid (8th order central differencing + WENO) Larsson & Gustavsson, 2008
  - High order central difference + Wavelet sensor based artificial dissipation  
ADPDIS3D (Yee & Sjogreen, 2007)
- Broad range of benchmark test cases
  - Noh problem (strong implosion)
  - Shu-Osher problem
  - Interaction of vorticity/Entropy wave with normal shock
  - Taylor-Green problem
  - Compressible isotropic turbulence

*Details in paper by Johnsen et al. to appear in J. Comput. Phys, 2009*



# Numerical Algorithm Assessment: Compressible Isotropic Turbulence (with eddy shocklets)

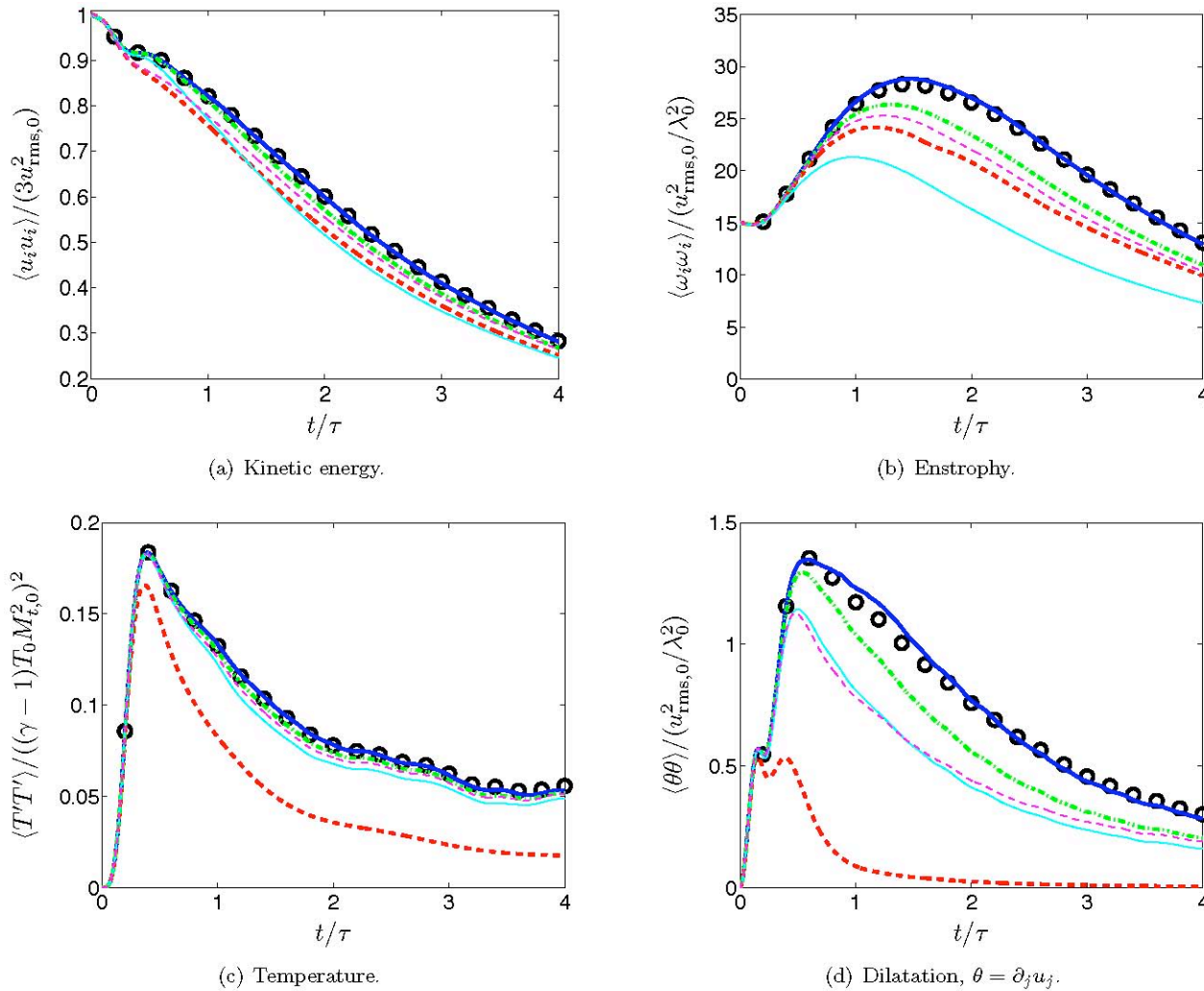
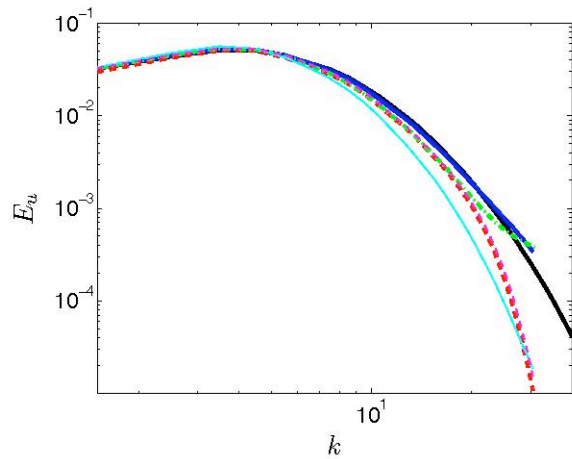


Figure 11: Temporal evolution of the variance of different quantities for the isotropic turbulence problem on  $64^3$  grid. The reference is the solution on a  $256^3$  grid spectrally filtered to a  $64^3$  grid (circles).

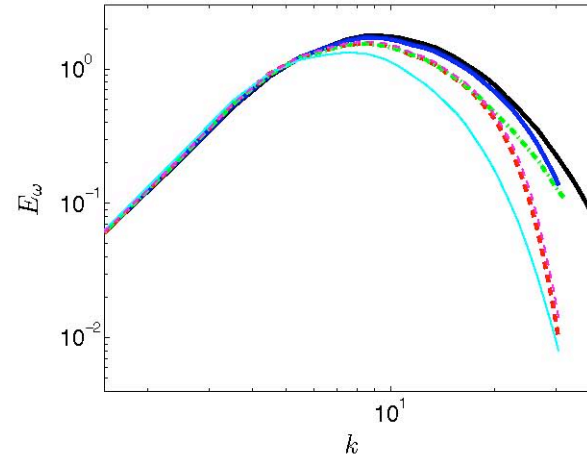
# Numerical Algorithm Assessment: Compressible Isotropic Turbulence (with eddy shocklets)



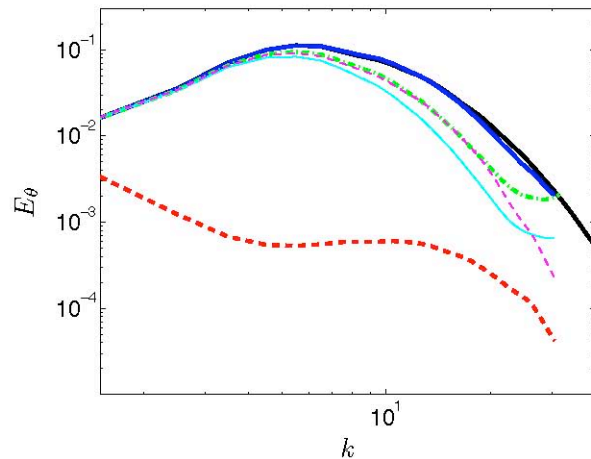
## Algorithm comparison



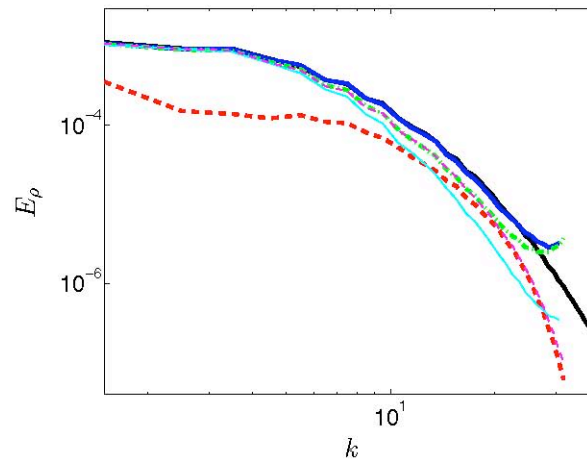
(a) Velocity.



(b) Vorticity.



(c) Dilatation.



(d) Density.

Code	Color	Line style
Reference	black	varying
<i>Stan</i>	red	dashed
<i>Stan-I</i>	magenta	dashed (thin)
<i>Hybrid</i>	blue	solid
<i>WENO</i>	cyan	solid (thin)
<i>ADPDIS3D</i>	green	dashed-dotted
<i>Shock Fit</i>	black	dotted

Table 2: Color and line legend for the plots.

Figure 13: Spectra at  $t/\tau = 4$  for the isotropic turbulence problem on a  $64^3$  grid. The reference is the solution on a  $256^3$  grid.



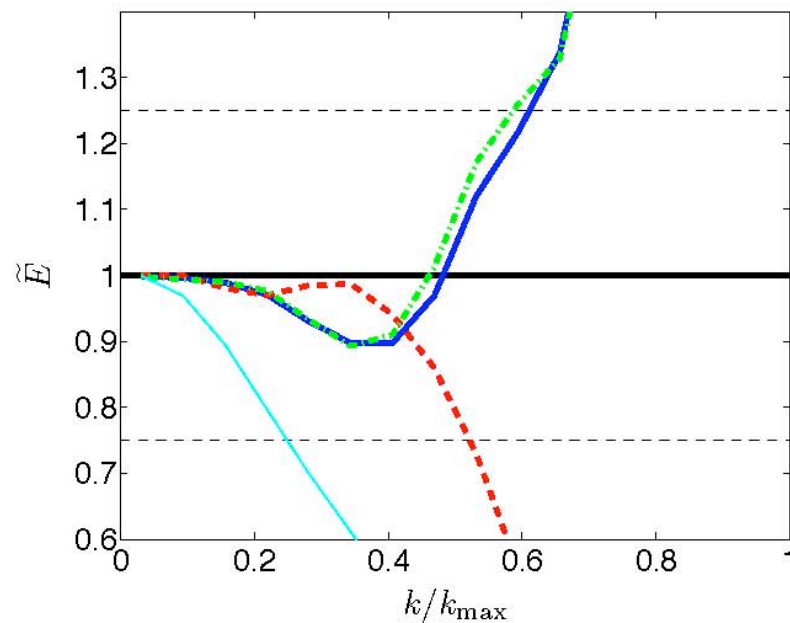
# Numerical Algorithm Assessment: Effective Bandwidth

- Use compensated spectra  $E(k)/E_{\text{exact}}(k)$  to define effective bandwidth

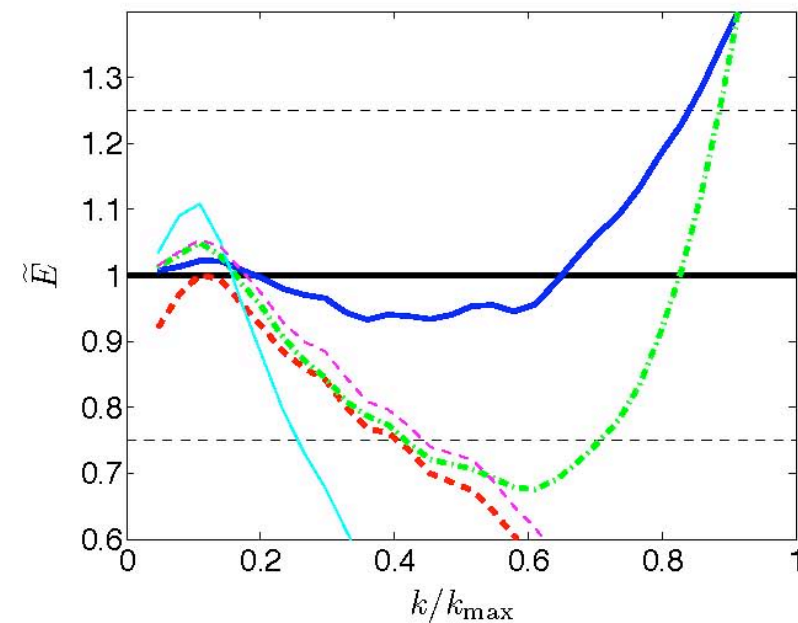
Code	Color	Line style
Reference	black	varying
<i>Stan</i>	red	dashed
<i>Stan-I</i>	magenta	dashed (thin)
<i>Hybrid</i>	blue	solid
<i>WENO</i>	cyan	solid (thin)
<i>ADPDIS3D</i>	green	dashed-dotted
<i>Shock Fit</i>	black	dotted

## Algorithm comparison

Table 2: Color and line legend for the plots.



(a) Taylor-Green vortex at  $t = 5$ .



(b) Isotropic turbulence at  $t/\tau = 4$ .

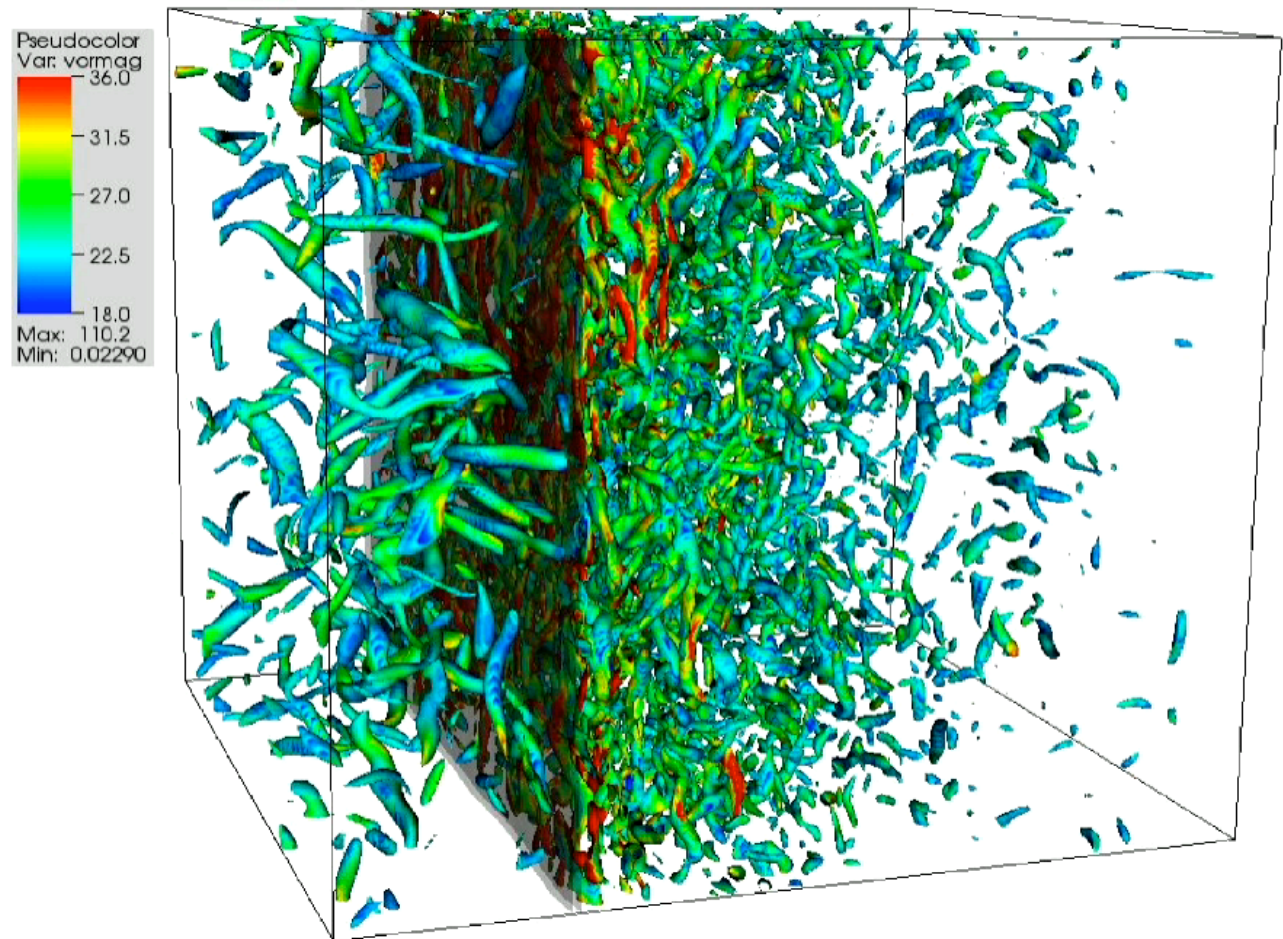
Figure 15: Compensated spectra on a  $64^3$  grid. As a reminder, *Stan* is dashed red, *Stan-I* is dashed magenta (thin), *Hybrid* is solid blue, *WENO* is solid cyan (thin), and *ADPDIS3D* is dash-dotted green.





# Canonical shock/turbulence interaction

- Isotropic turbulence passing through a normal shock in a perfect gas
- Isolates the core interaction between turbulence and shock



Eddies visualized by Q-criterion, colored by vorticity. Shock visualized by dilatation contour.



## Some past work on canonical shock/turbulence interaction

- Theoretical
  - Ribner (NACA 1953, NACA 1954, AIAA J 1987), Moore (NACA 1954): linear interaction analysis including Rankine-Hugoniot relations
  - Lele (PhysFI 1992): turbulent shock relations using RDT, predicted modified mean shock jumps
- Experimental: wind tunnels
  - Barre et al (AIAA J 1996): grid turbulence passing through a Mach 3 shock, hot-wire and LDV measurements
- Experimental: shock tubes
  - Hesselink & Sturtevant (JFM 1988): found “peaked” and “rounded” instantaneous pressure profiles in random medium
- Computational (DNS):
  - Lee et al (JFM 1993): resolved viscous shock structure,  $M = 1.2$ ,  $Re_\lambda = 20$ ,  $M_t = 0.1$ , found modified instantaneous profiles of dilatation
  - Lee et al (JFM 1997): captured shock up to  $M = 3$ , found good agreement with Ribner’s linear theory
  - Mahesh et al (JFM 1997): influence of entropy fluctuations
  - Ducros et al (2000): LES; Jamme et al (FTC 2002)



## Post-shock Kolmogorov length scale: Scaling Estimate

- From the (incompressible) definition of the Kolmogorov length scale

$$\eta = \left( \frac{\nu^3}{\varepsilon/\rho} \right)^{1/4} \sim \left( \frac{\mu^2}{\rho^2 \omega^2} \right)^{1/4} \sim \frac{T^{3/8}}{\rho^{1/2} \omega^{1/2}}$$

where we used

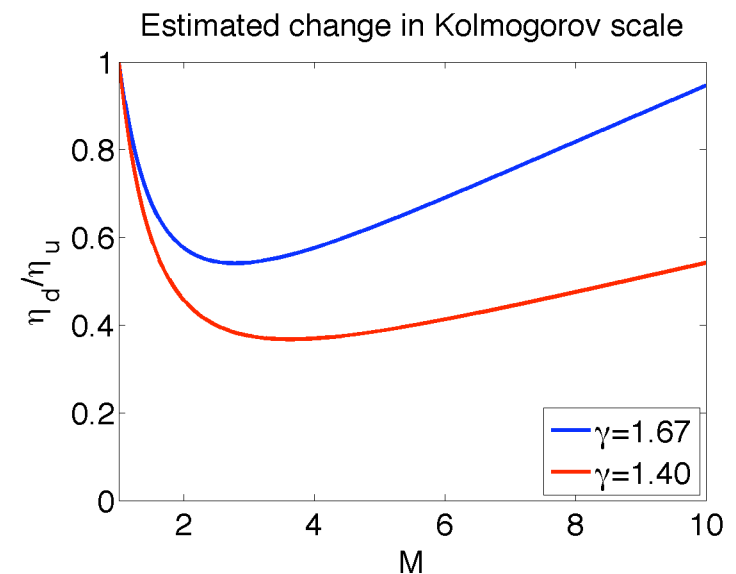
$$\varepsilon \approx \mu \omega^2, \quad \nu = \mu/\rho, \quad \mu \sim T^{3/4}$$

- Across the shock, assume that the vorticity changes as the density ratio; gives

$$\frac{\eta_d}{\eta_u} \approx \left( \frac{T_d}{T_u} \right)^{3/8} \left( \frac{\rho_d}{\rho_u} \right)^{-1}$$

in terms of Rankine-Hugoniot relations

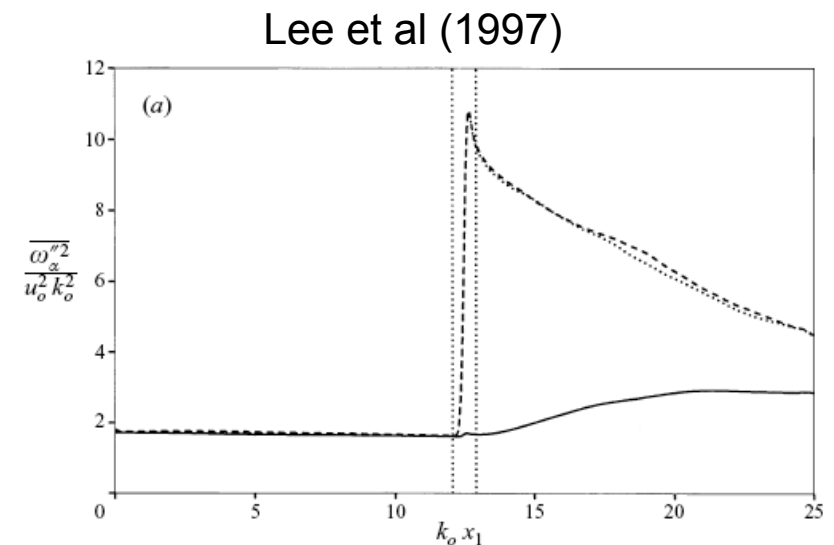
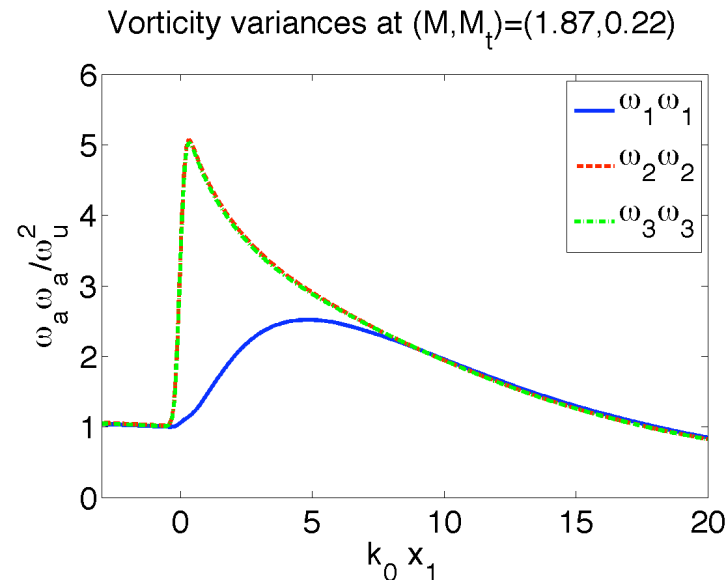
Post-shock Kolmogorov Scale determines the resolution requirement for turbulence





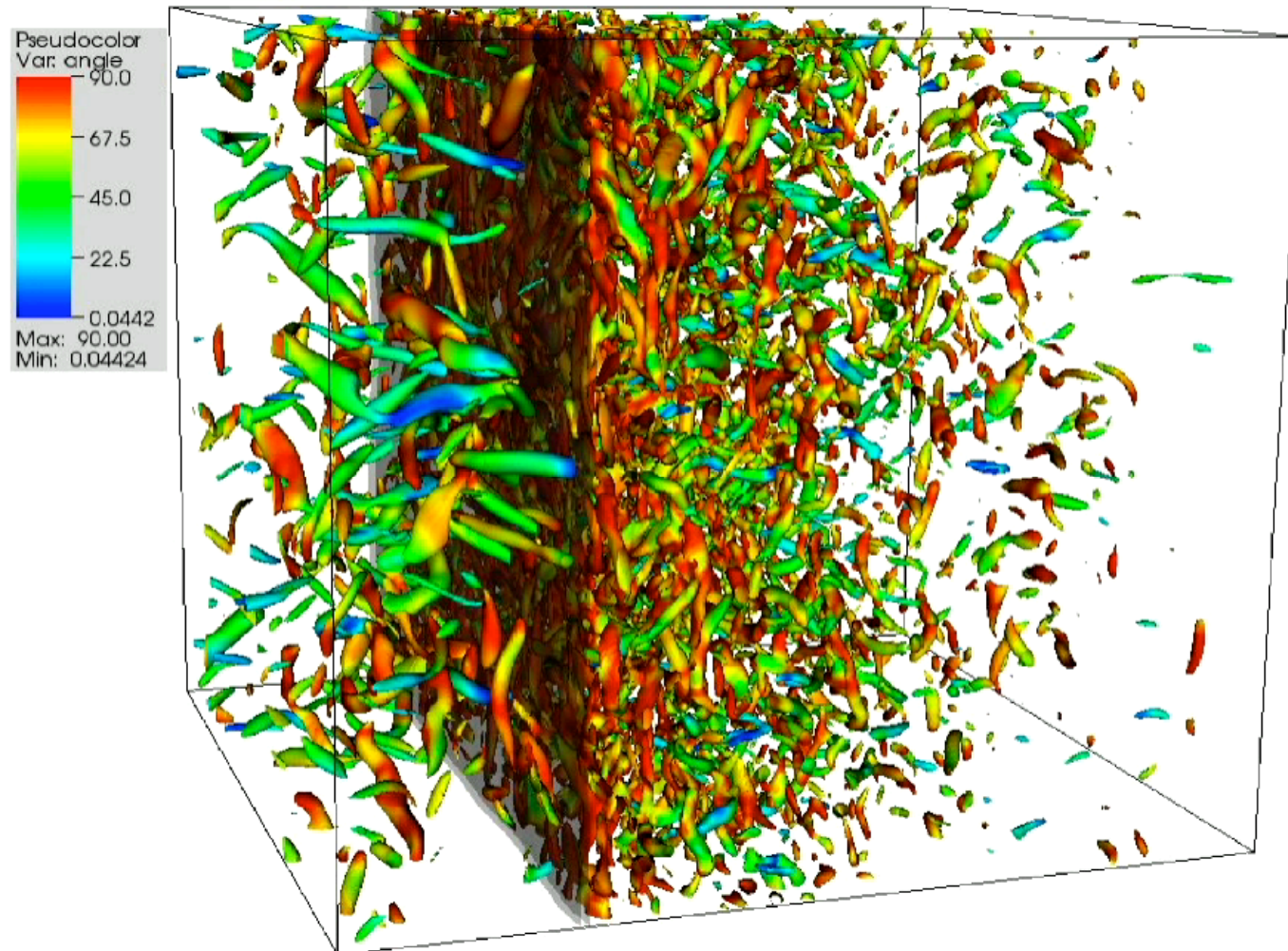
## Difference compared to previous studies

- Past studies:  $Re_\lambda = 20$  on  $129 \times 64^2$  grids (Lee et al, 1993, 1997)
- Present:  $Re_\lambda = 40$  on  $1040 \times 384^2$  grids;  $Re_\lambda = 60$  on  $1675 \times 512^2$  grids
- Return to local (small-scale) isotropy in present DNS, but not in past work
- Vorticity redistribution is a nonlinear process
  - Past work did not fully capture this due to lack of grid resolution
  - Has led to persistent misunderstanding of vorticity evolution at shock





# Anisotropic post-shock turbulence

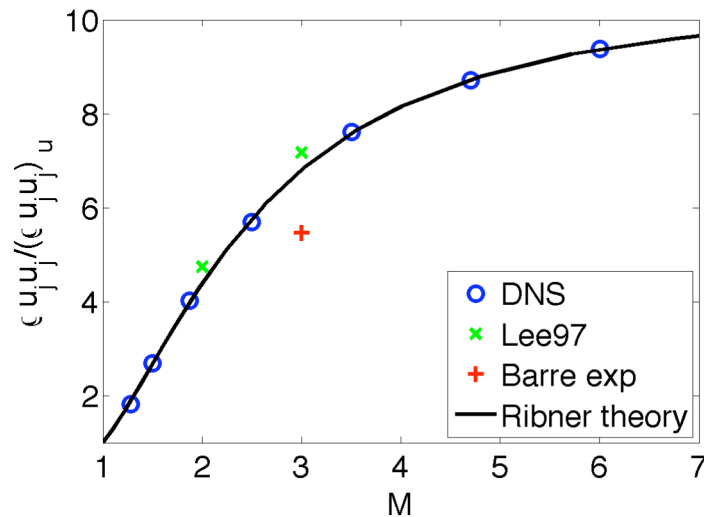


Eddies visualized by Q-criterion, colored by angle between vorticity vector and x-axis.  
Shock visualized by dilatation contour.

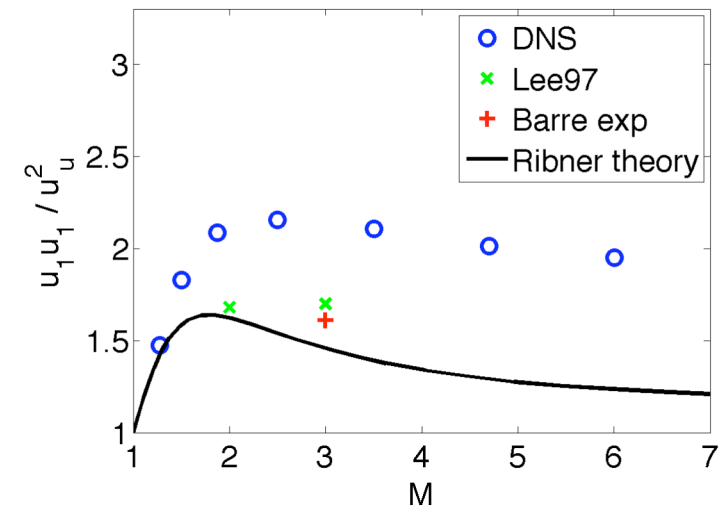


# Reynolds stresses: amplification across the shock vs Mach number

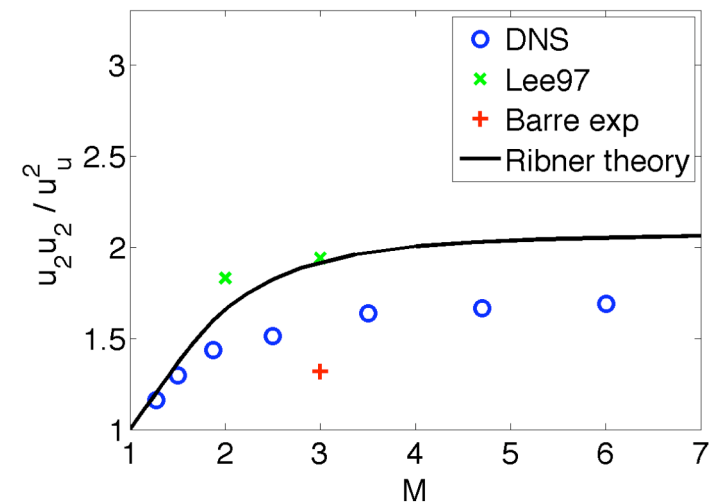
Amplification of kinetic energy at  $M_t=0.22$



Amplification of  $u_1 u_1$  at  $M_t=0.22$



Amplification of  $u_2 u_2$  at  $M_t=0.22$

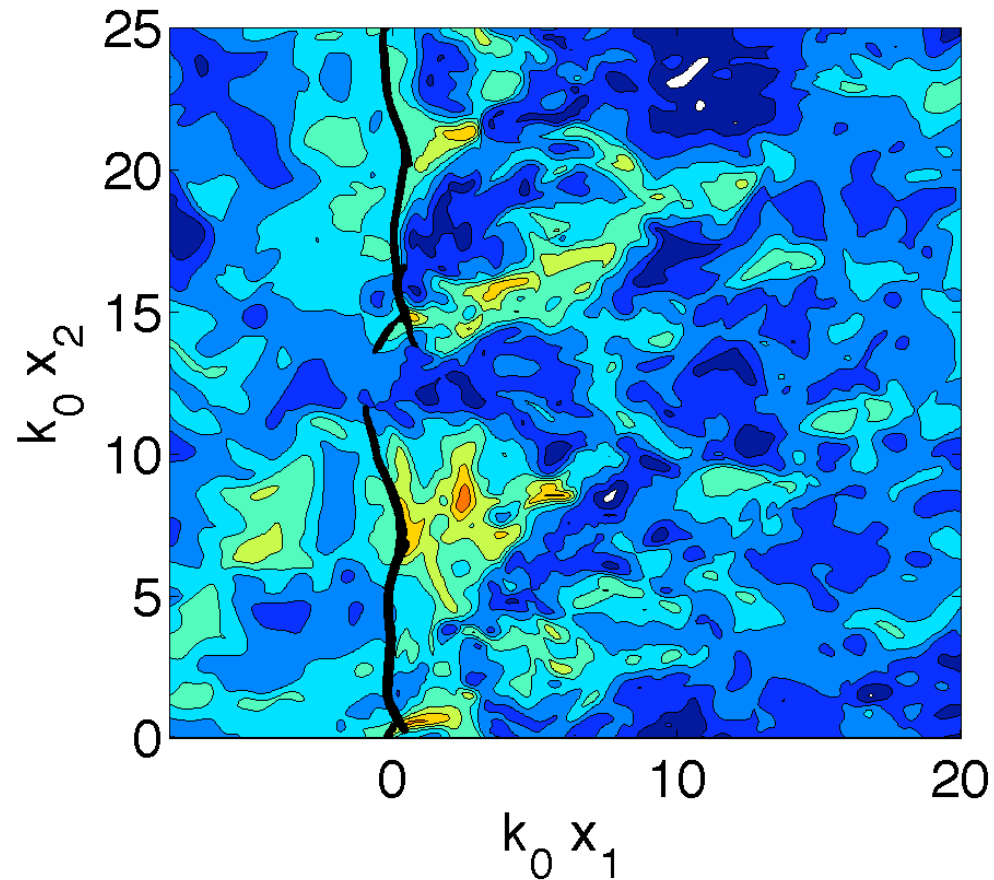


- Amplification of kinetic energy agrees very well with Ribner's linear theory (and Lee97)
  - Kinetic energy must be governed largely by linear processes
- DNS (and experiment) yields much larger anisotropy
  - Nonlinear redistribution processes must be important



## Complex shock structure in nonlinear regime

x-momentum at  $(M, M_t) = (1.50, 0.38)$

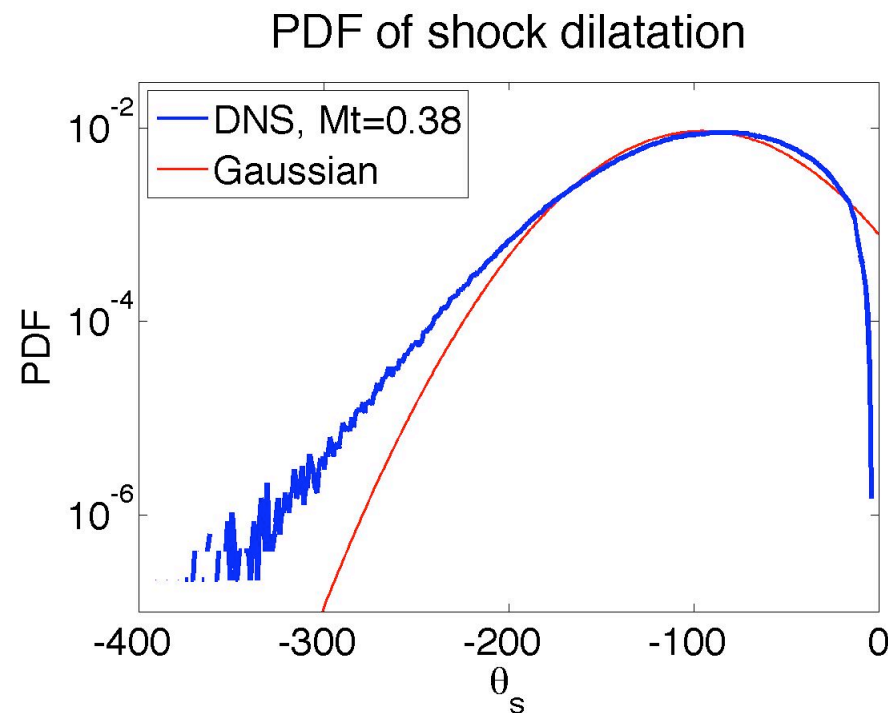


- Shock no longer a simple object
  - Qualitative agreement with Zank et al (PhysFI 2002)



## Extracting “extreme” interaction events

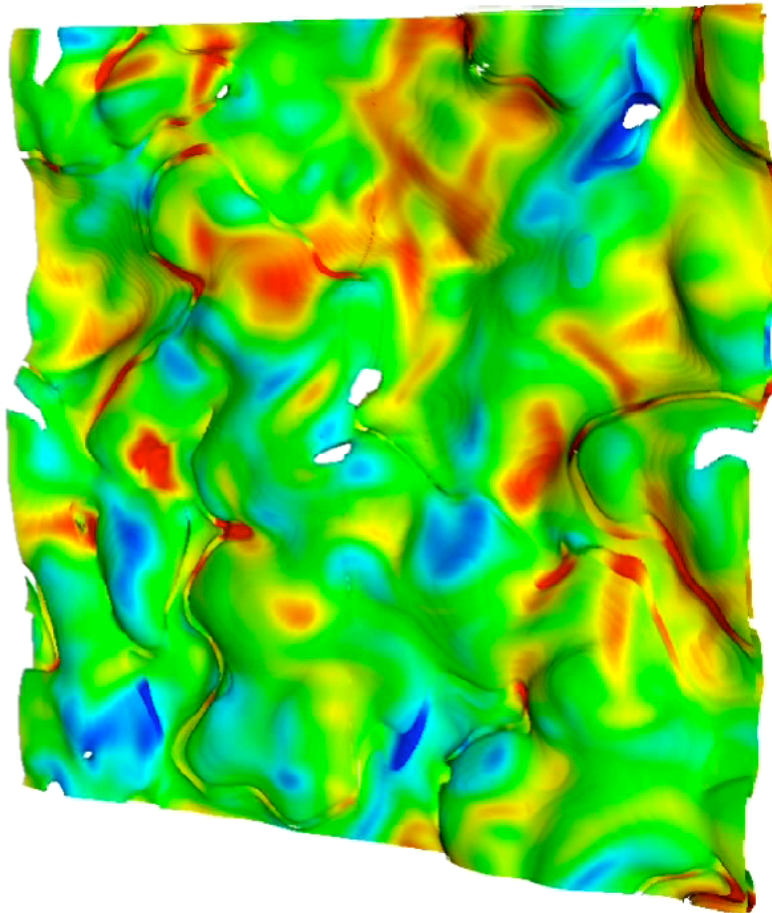
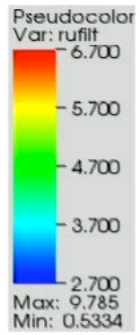
- Define shock-dilatation as  $\theta_{\text{shock}}(y, z) = \min_x \theta(x, y, z)$
- Study instantaneous profiles through the shock at extreme values of shock-dilatation:
  - “Strong” interactions: most compressive events
  - “Weak” interactions: least compressive events



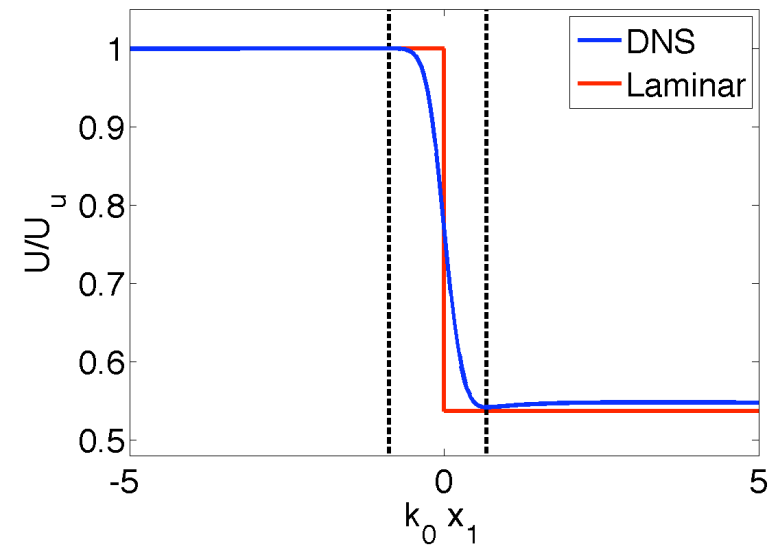




# Unsteady shock motion



Mean velocity at  $(M, M_t) = (1.50, 0.22)$



Shock visualized by dilatation contour, colored by streamwise momentum.

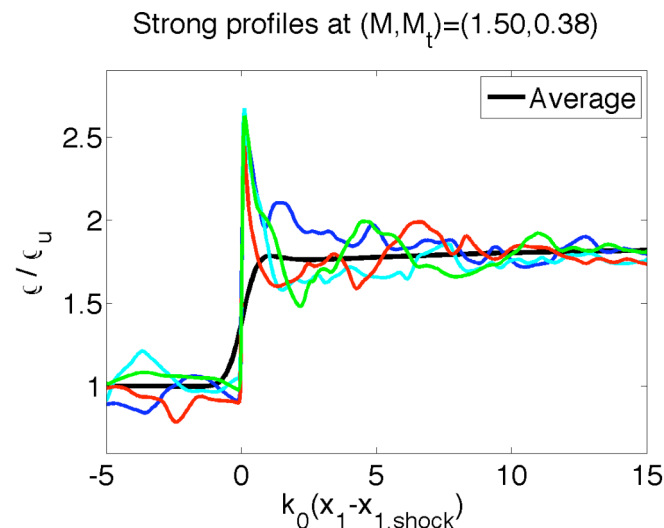
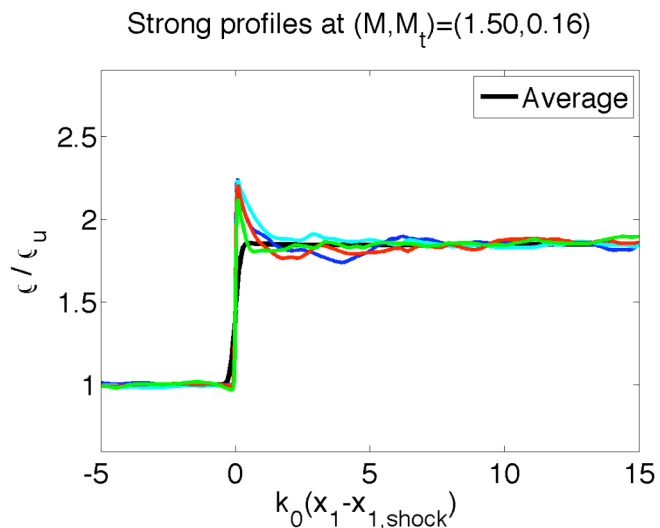


# “Strong” events: over-compression and post-shock expansion

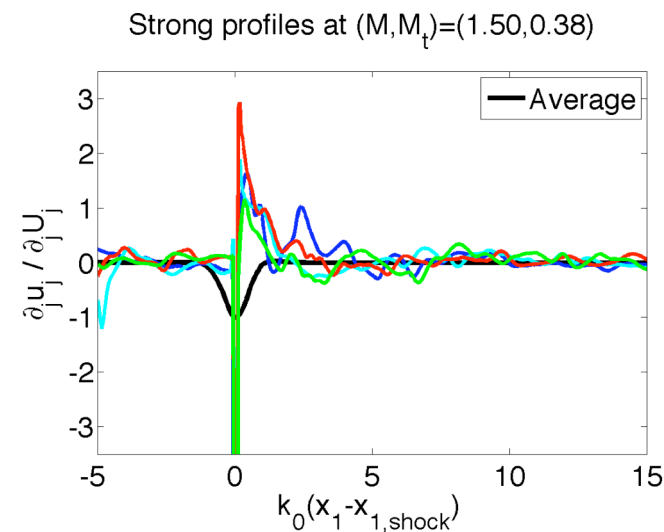
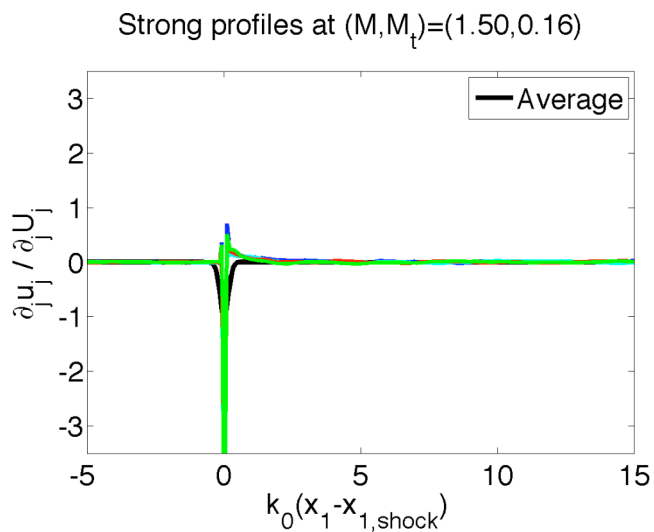
Low  $M_t$

High  $M_t$

Density



Dilatation





# “Weak” events: smooth profiles in nonlinear (high $M_t$ ) regime

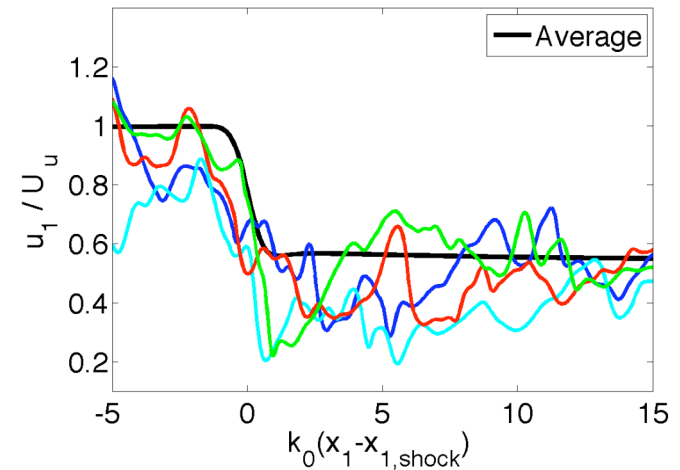
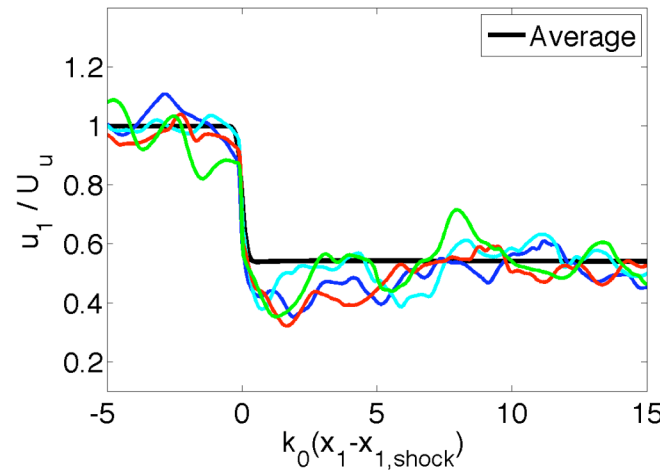
Low  $M_t$

High  $M_t$

Weak profiles at  $(M, M_t) = (1.50, 0.16)$

Weak profiles at  $(M, M_t) = (1.50, 0.38)$

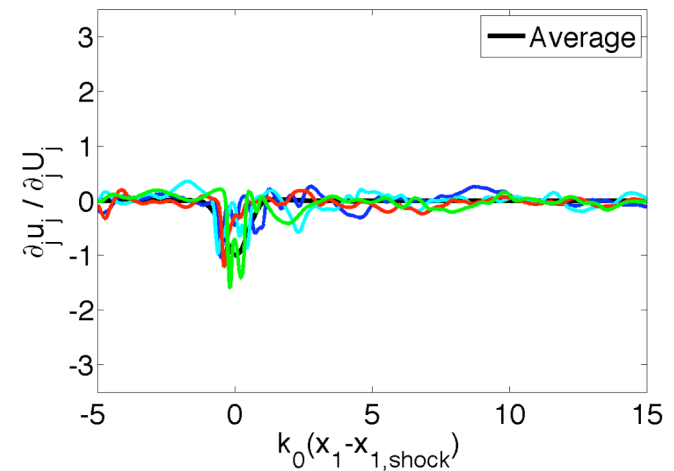
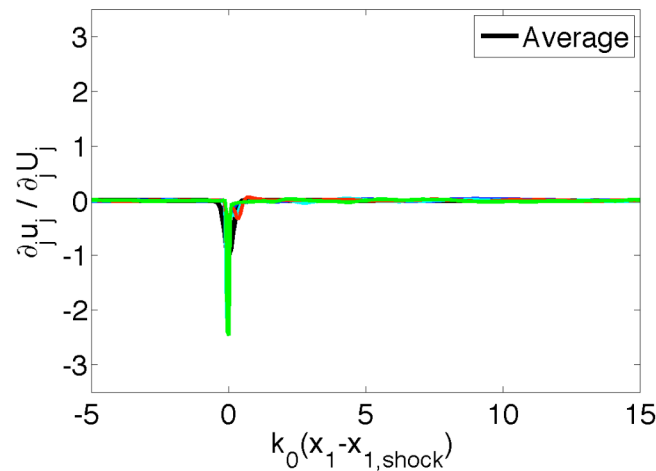
Velocity



Weak profiles at  $(M, M_t) = (1.50, 0.16)$

Weak profiles at  $(M, M_t) = (1.50, 0.38)$

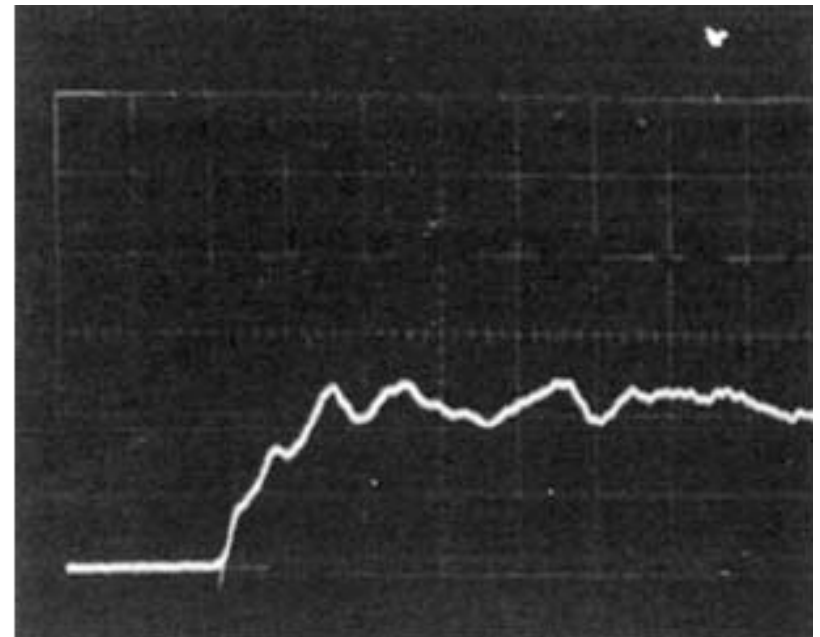
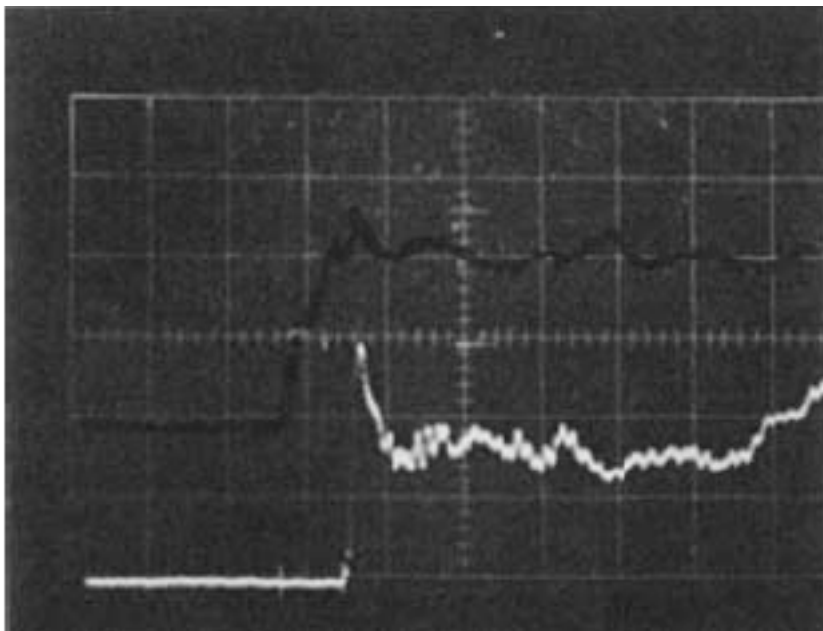
Dilatation





## Hesselink & Sturtevant (JFM 1988)

- “Peaked” and “rounded” pressure traces for shock propagating through a random mixture of He and R12 at  $M = 1.007 - 1.1$





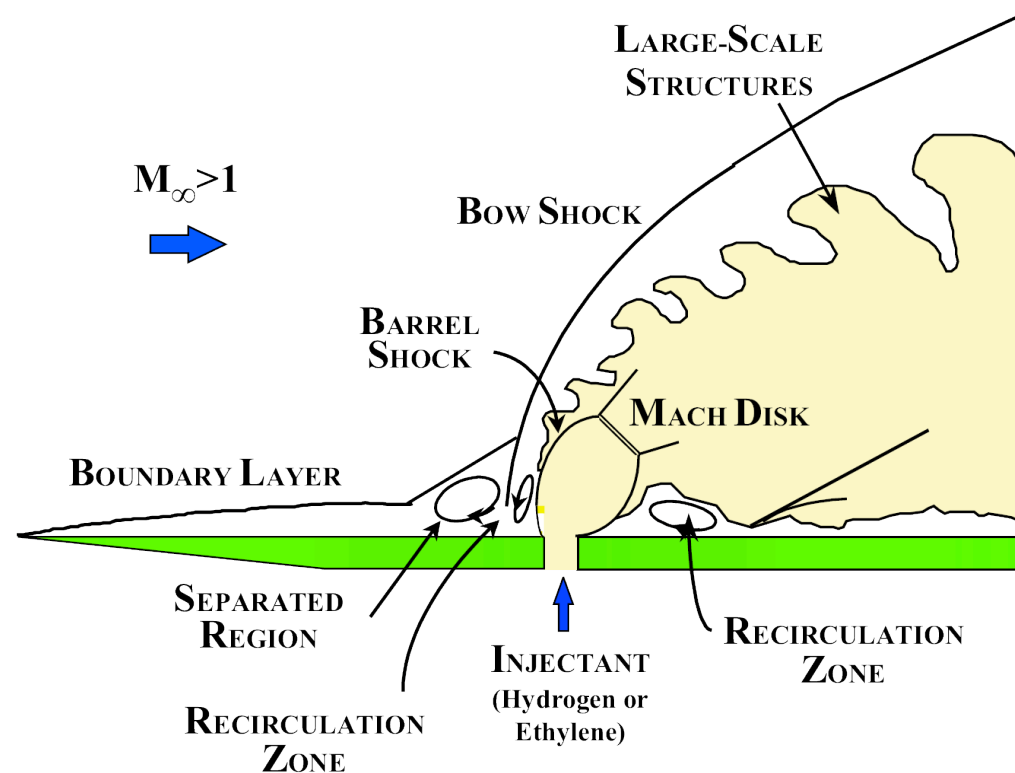
## Summary -- what we've learned so far

---

- Direct numerical simulations of canonical shock/turbulence interaction that truly resolves the post-shock turbulence
- Disagreements with linear theory and past computations:
  - Vorticity components return to isotropy
  - Qualitatively different, and much larger, stress anisotropy
- Agreement with linear theory and past computations:
  - Kinetic energy amplification -- governed by linear processes
- Instantaneous shock/turbulence interaction dynamics:
  - Shock weakens (to the point of disappearing!) locally -- smooth profiles through shock in nonlinear regime
  - “Strong” shock events with up to twice the average compression



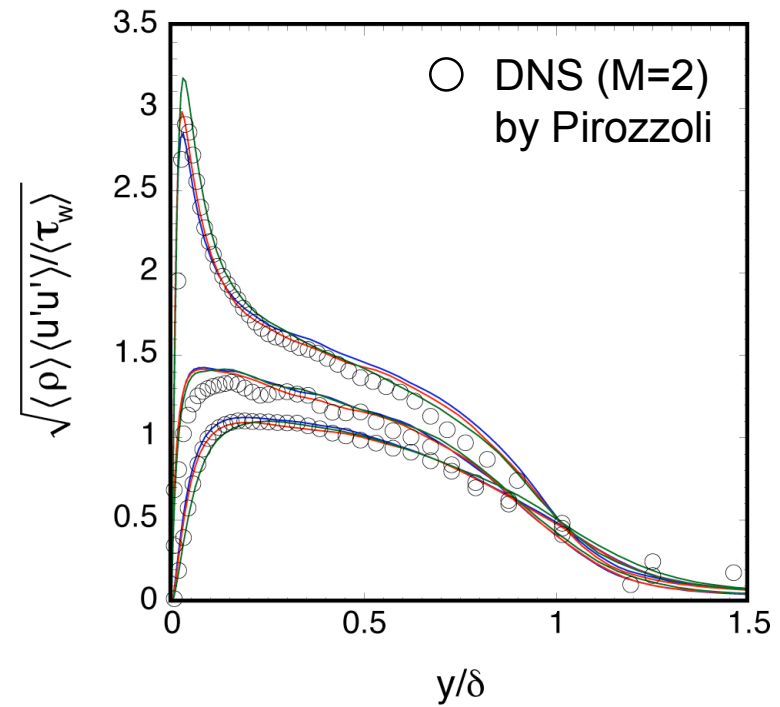
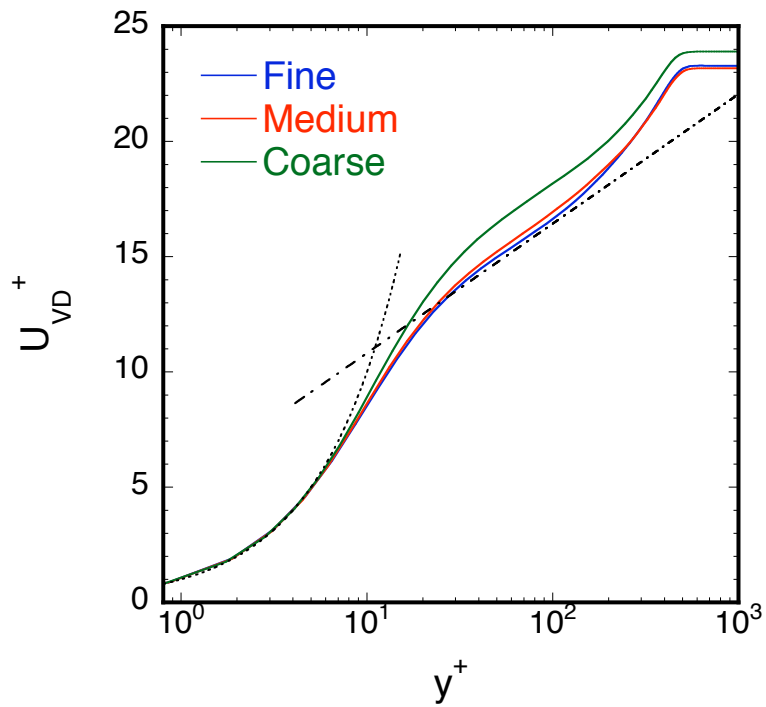
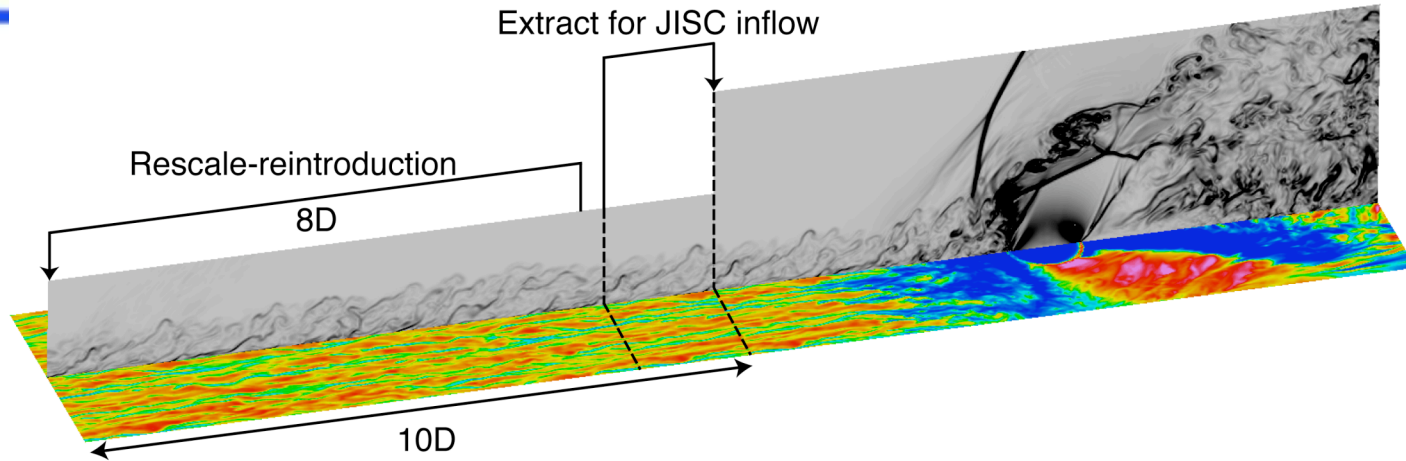
# A more complex flow with shock-turbulence interaction



Expt. By Santiago & Dutton, 1997, JPP; Ben Yakar et al. 2006, PoF  
LES by Kawai & Lele, AIAA 2009-3995



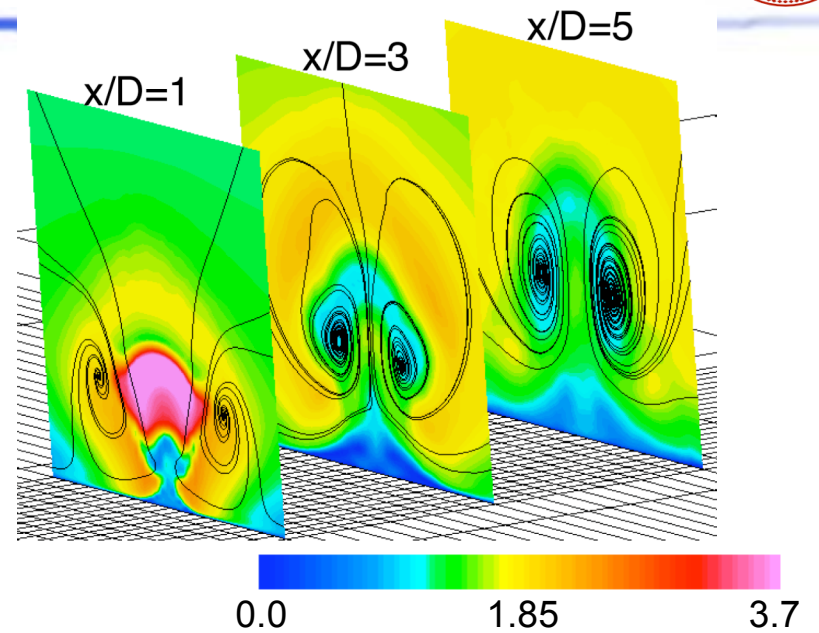
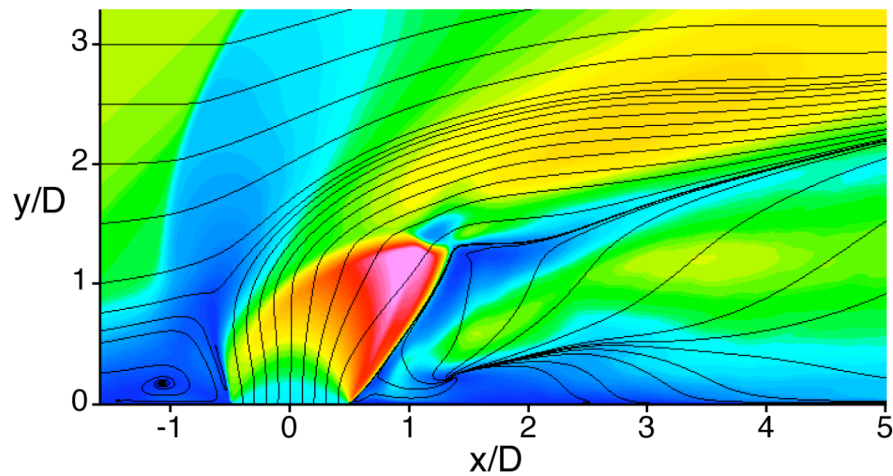
# JISC with Turbulent Crossflow



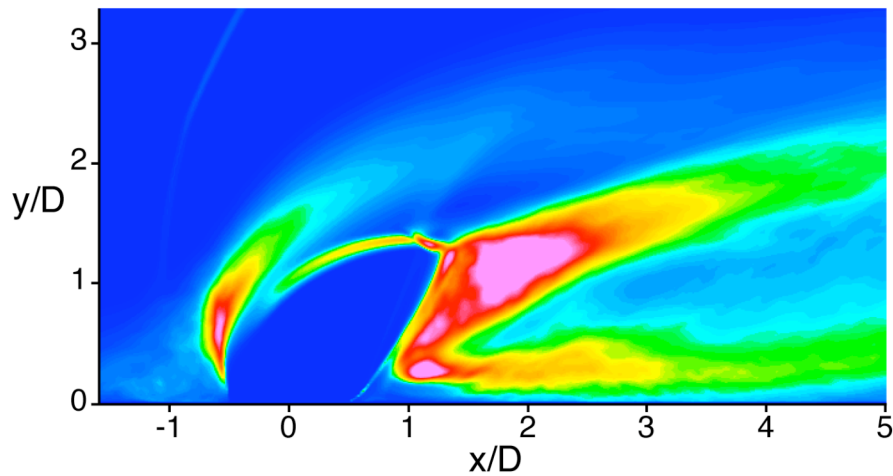


# Time-Averaged Flowfields

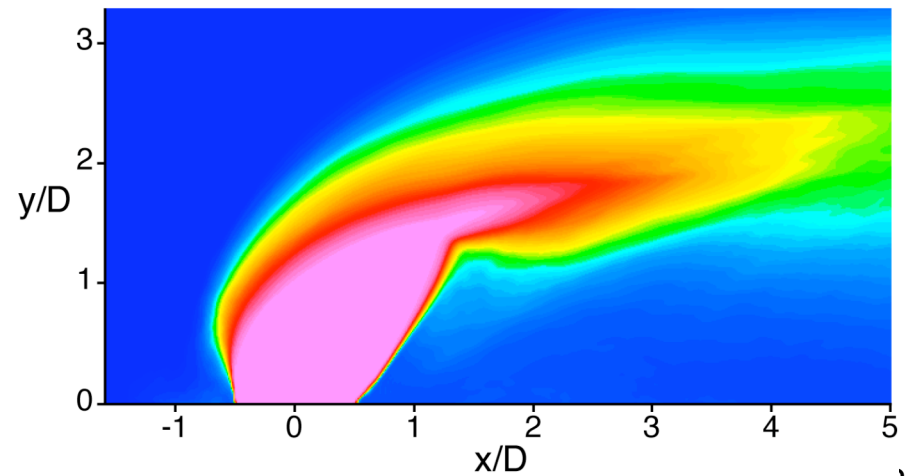
## Mach number



## Turbulent kinetic energy



## Jet fluid (passive scalar)





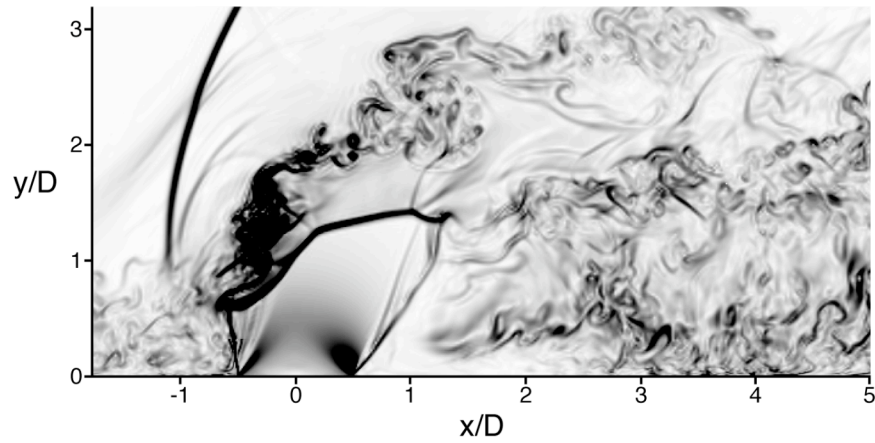


# Animation

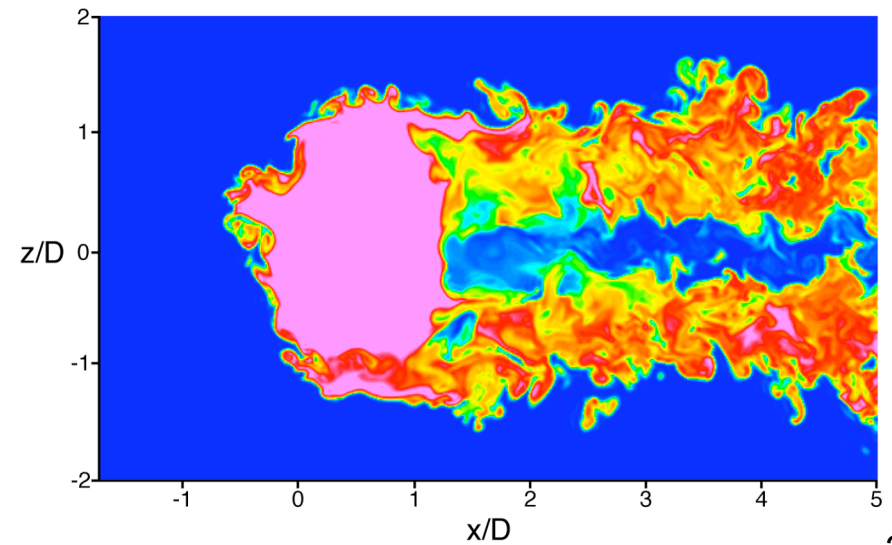
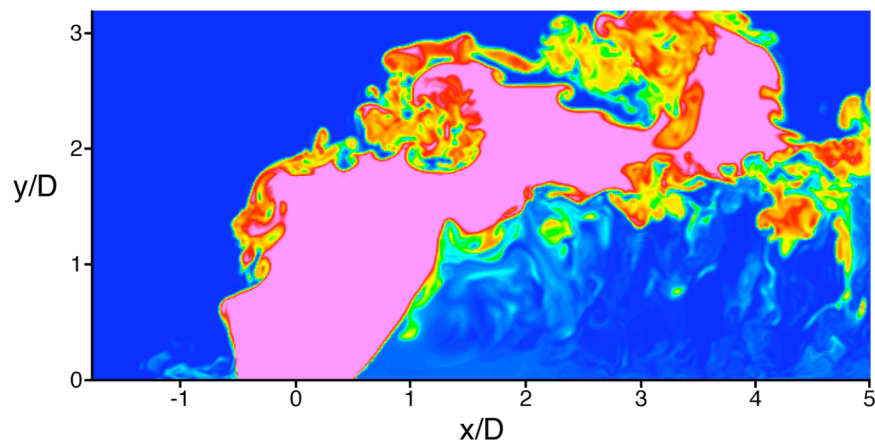
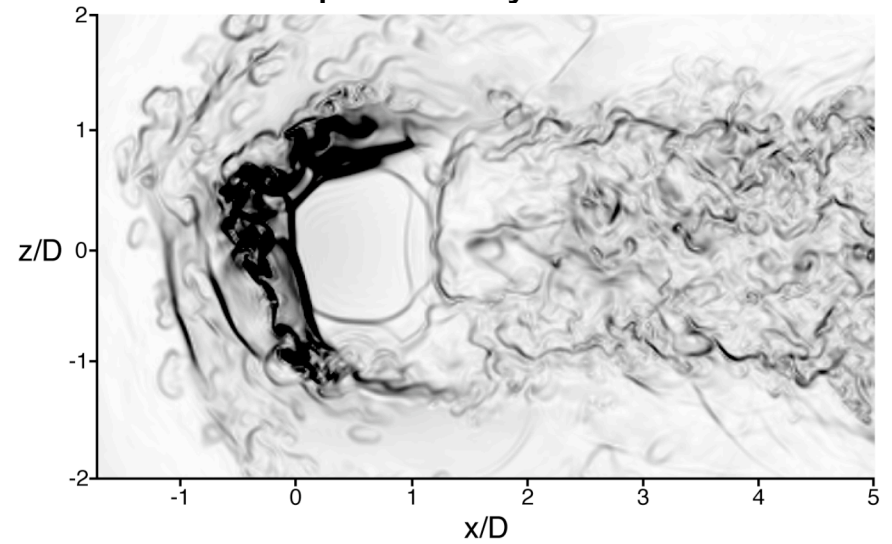
Norm of density gradient

Jet fluid (scalar)

Sideview at  $z/D=0$

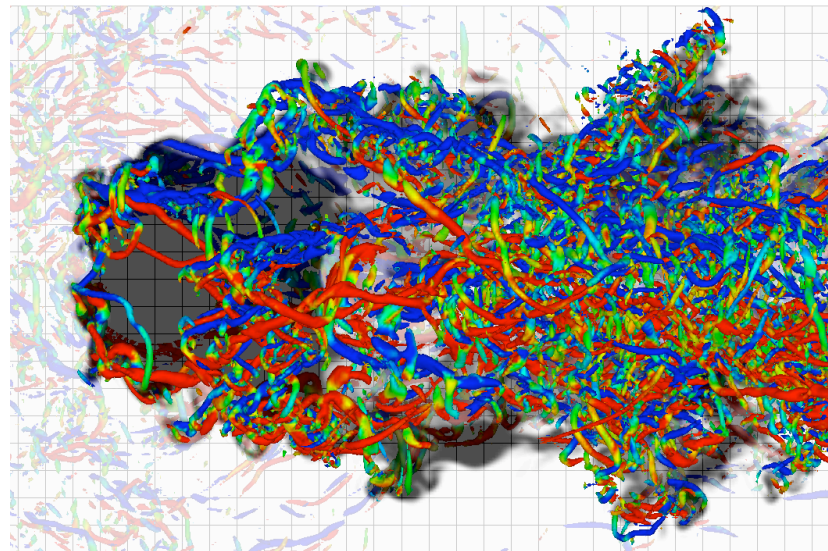
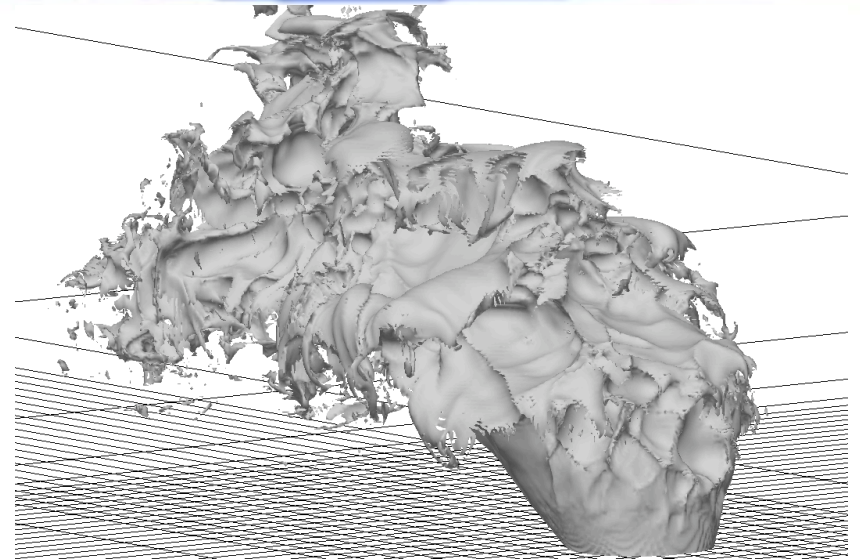
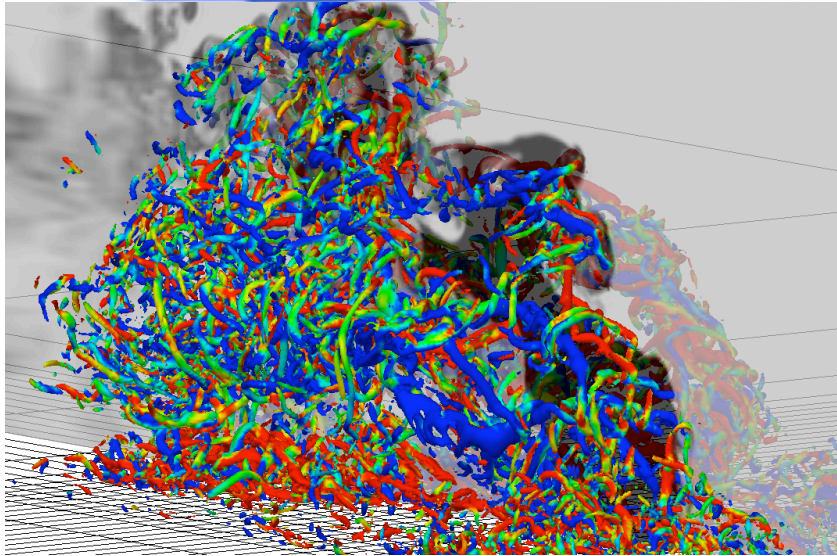


Topview at  $y/D=1$





# 3D Flow Structures





# Closure ...

---

- Revisit to canonical shock-turbulence interaction (DNS)
- Other studies (DNS/LES) of shock-turbulent BL interaction, compression ramp..
- More complex engineering flows involving injection, mixing and combustion in high-speed flows

These are good laboratories for improving our understanding of turbulence behavior, its engineering modeling (RANS & LES), and for studies of aeroacoustics and flow control.

What will the next 50 years bring ?

Exciting science and collaborative research for sure. Happy 50 ! And Thanks.



Additional slides



## Summary -- new and unanticipated questions

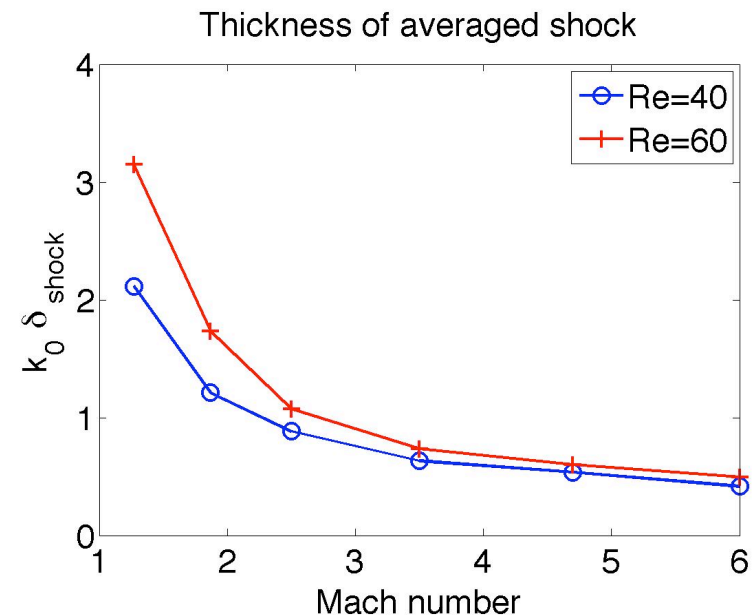
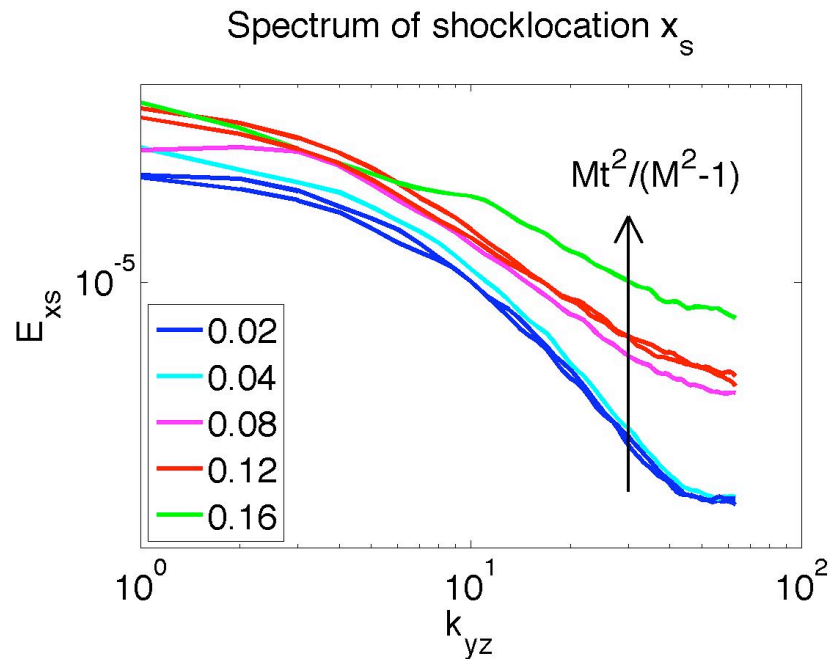
---

- Effect of turbulent Mach number on amplification of turbulence -- viscous effects make conclusions difficult
  - New  $Re_\lambda = 60$  runs should clarify this
- Why is peak shock-motion at larger length scales than peak turbulence energy?
  - Theory suggests shock-motion “enslaved” to incoming turbulence
  - Need runs with higher  $k_0$  to allow for deeper analysis
- What is the relation between shock-motion and Reynolds number? Or is the shock-motion dependent on some unknown parameter that was not matched?
- What is the turbulence structure behind the shock -- statistics show elongation in streamwise direction, but what does it look like?
  - More, and more intelligent, visualization needed



# Unsteady shock motion

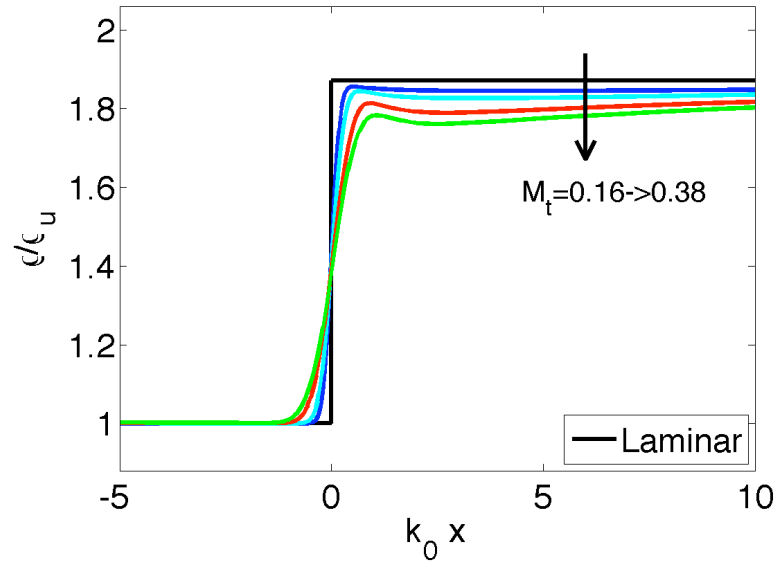
- Interesting finding: spectrum of shock motion seems to scale with ratio of turbulent pressure fluctuation to shock-induced pressure jump:  $M_t^2 / (M^2 - 1)$
- Puzzle #1: turbulence has peak energy at  $k=4$ , but shock motion has peak at  $k=1-2$
- Puzzle #2: thickness of averaged shock increases with  $Re$



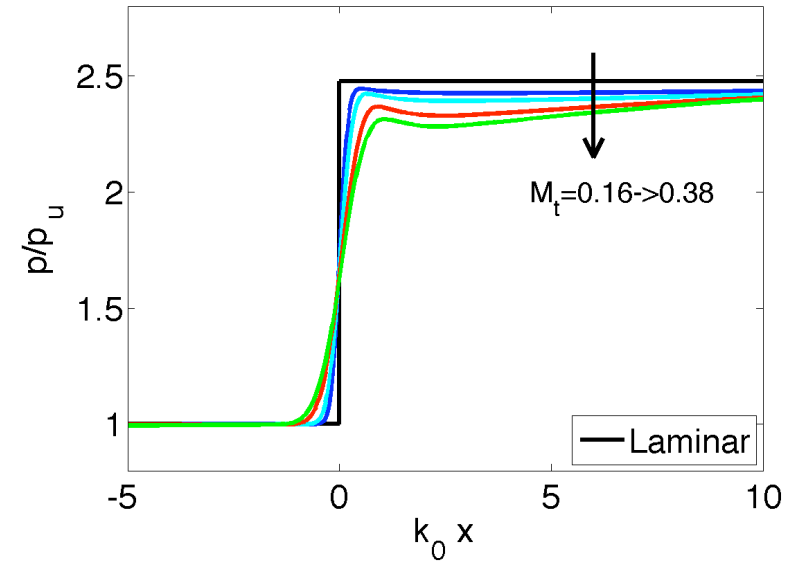


# Mean profiles

Mean density at M=1.50



Mean pressure at M=1.50



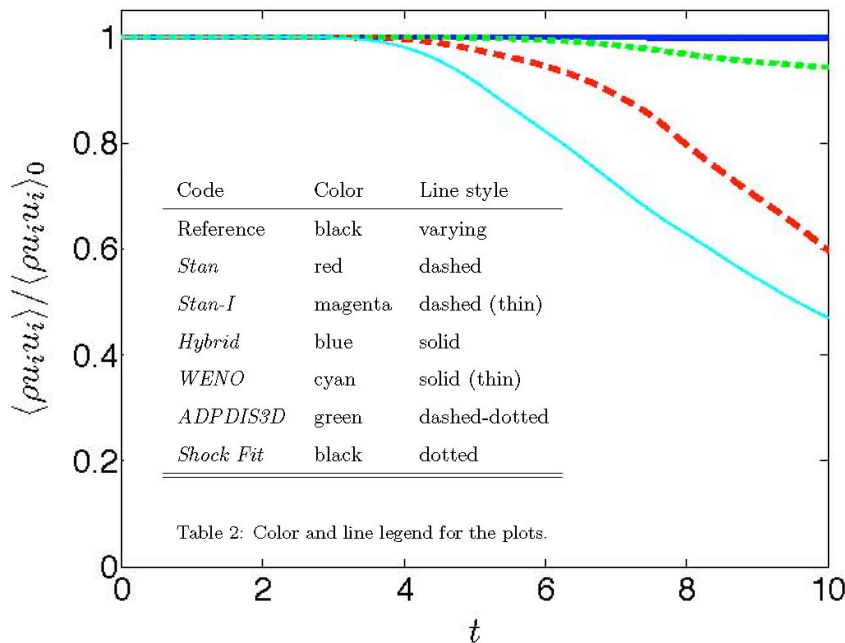
- Turbulence modifies the Rankine-Hugoniot relations (Lele, PhysFI 1992):

$$\begin{aligned} \overline{\rho \tilde{u}} &= \text{const} \\ \overline{\rho \tilde{u}^2} + \overline{p} + \overline{\rho \widetilde{u'' u''}} &= \text{const} \\ \overline{\rho h_0 \tilde{u}} + \overline{\rho \widetilde{h_0'' u''}} &= \text{const} \end{aligned}$$

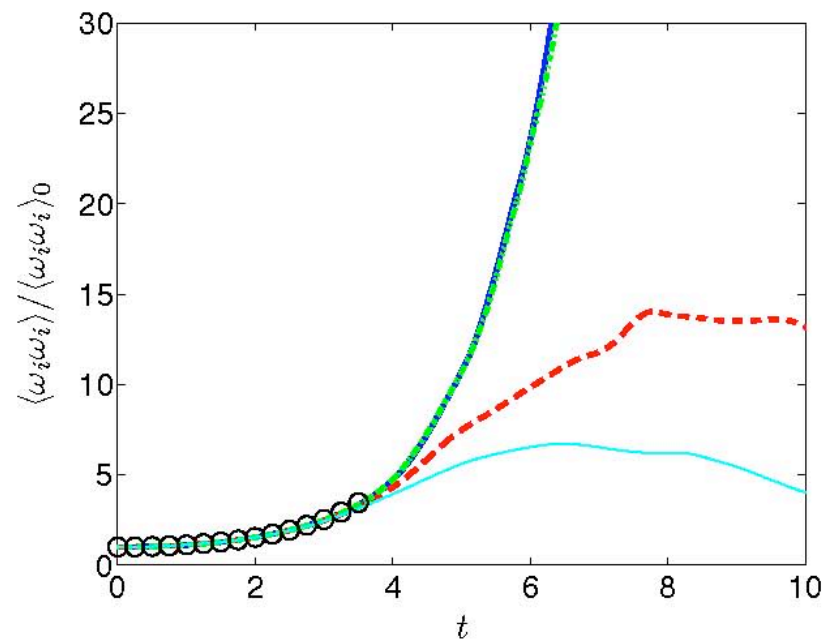
# Numerical Algorithm Assessment: Taylor-Green Problem



- Inviscid problem with strong vortex-stretching  
Kinetic energy conservation



(a) Kinetic energy.



(b) Enstrophy. The semi-analytical result of Brachet *et al.* (1983) are the black symbols.

Figure 1: Mean quantities for the Taylor-Green vortex on a  $64^3$  grid. The zero subscript denotes the initial value.

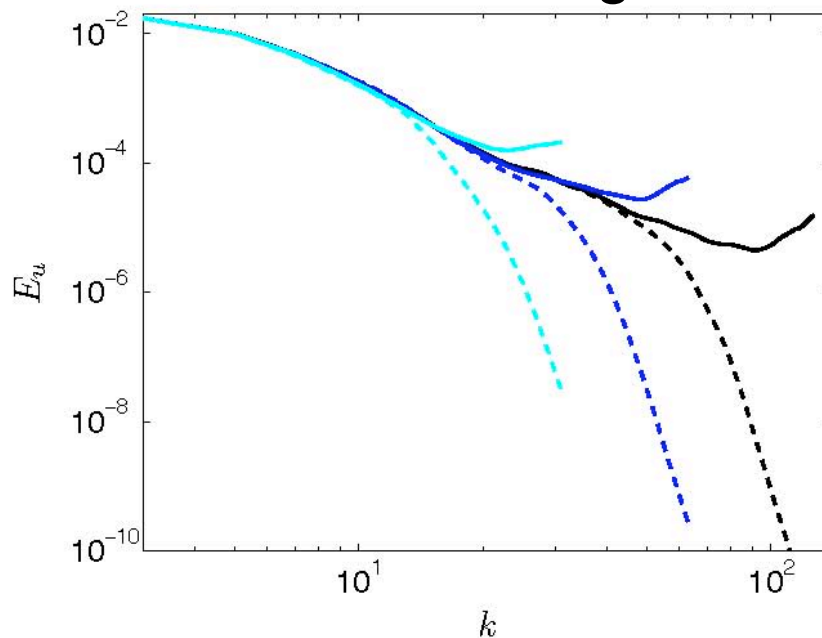


# Numerical Algorithm Assessment: Taylor-Green Problem

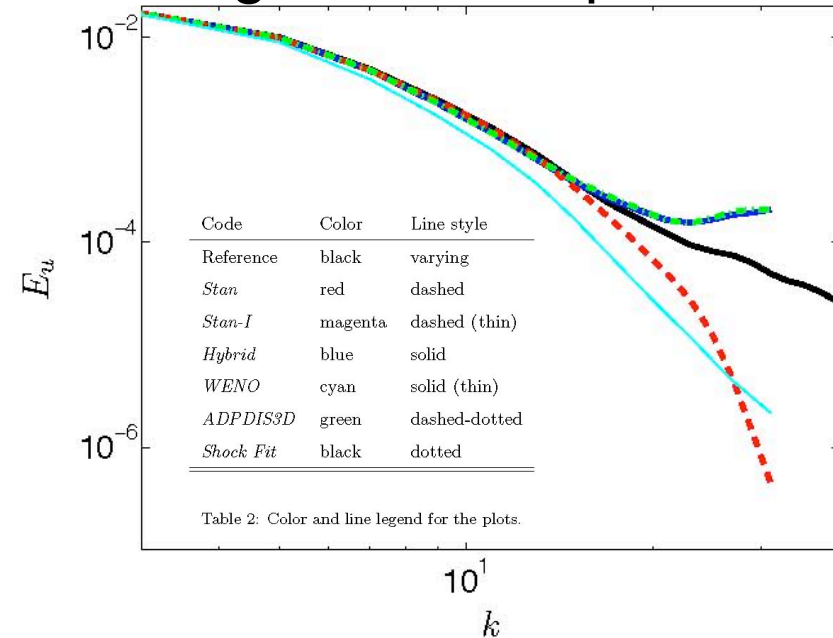


- Kinetic Energy Spectrum

## Grid Convergence



## Algorithm comparison



(a) Convergence of the reference solution using the *Hybrid* code in standard mode (solid) and with eight-order accurate dissipation (dashed), on  $256^3$  (black),  $128^3$  (blue), and  $64^3$  (cyan).

(b) Comparison between the different schemes. The reference solution is that obtained on the  $256^3$  grid using the *Hybrid* code.

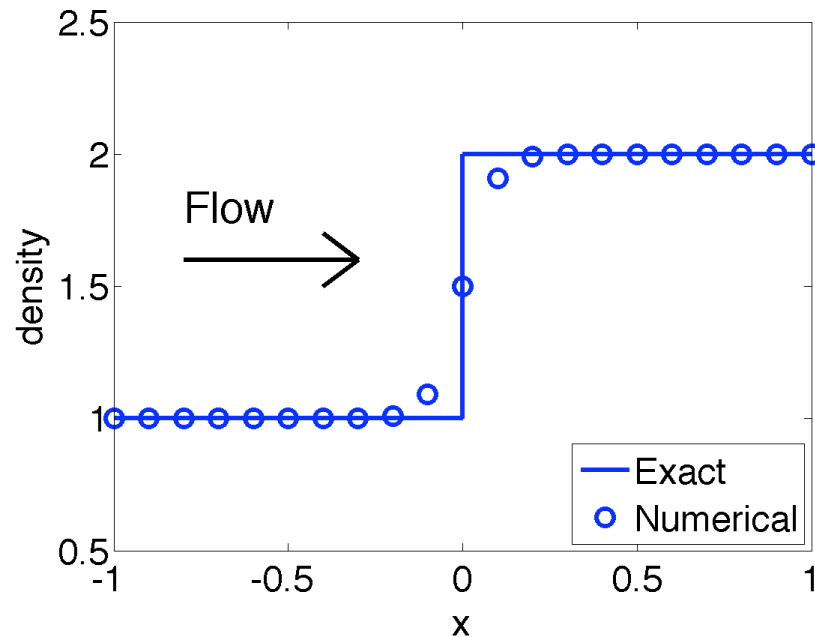
Figure 2: Velocity spectra for the Taylor-Green vortex on  $64^3$  grid at  $t = 5$ .



# Direct numerical simulation (DNS)

- Solve the Navier-Stokes equations without modeling
- Numerical challenges:
  - Shock-capturing requires numerical dissipation to smear shock
  - Numerical dissipation kills turbulence

Numerically captured shock



Temperature variance

